

Rarefied Gas Dynamics: Theory and Applications to Vacuum

Lecture 2: Direct Simulation Monte Carlo

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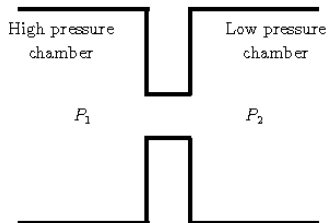
<http://fisica.ufpr.br/sharipov>

XLIV CBrAVIC, Short Course

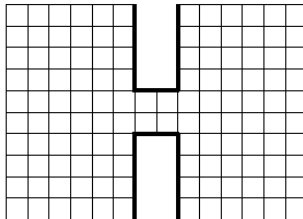
São Paulo, November 25, 2023

Outline

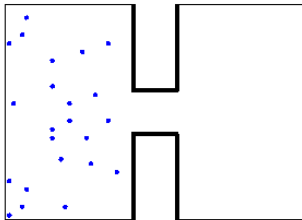
- Direct simulation Monte Carlo method.
- Intermolecular collisions
- Hard sphere vs. ab initio potential
- Orifice flow
- Modelling of vacuum pumps
- 3D flows
- Transient flows



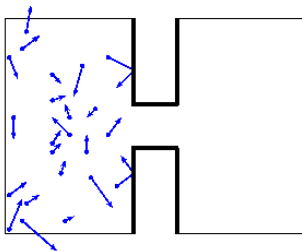
Example: Gas flow through a short tube.



Flow region is divided into cells



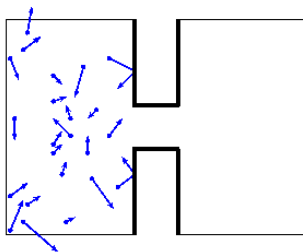
M model particles are considered.
Their positions \mathbf{r}_i and velocities \mathbf{v}_i are stored.



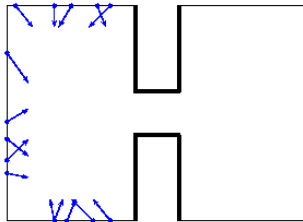
Time is advanced in steps Δt .

Free motion without collisions. New positions are calculated

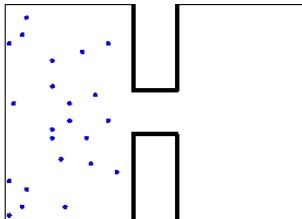
$$\mathbf{r}_{i,new} = \mathbf{r}_{i,old} + \mathbf{v}_i \Delta t \quad (1)$$



Gas-surface interaction is simulated.
Some particles are removed.



New particles are generated.



Intermolecular collisions are simulated. New velocities v_i are calculated
Macroscopic quantities are calculated.

All steps are repeated many times in order to reduce the statistical noise.

Diffuse scattering on solid surface

Normal component of velocity

$$f(v_n) = 2 \frac{v_n}{v_w^2} \exp\left(-\frac{v_n^2}{v_w^2}\right), \quad v_w = \sqrt{\frac{2k_B T_w}{m}} \quad 0 \leq v_n < \infty \quad (2)$$

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$$v_n = v_w \sqrt{-\ln R_f}, \quad R_f - \text{random fraction of } 1 \quad (3)$$

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Diffuse scattering on solid surface

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$$v_{t1} = v_t \cos \phi, \quad v_{t2} = v_t \sin \phi \quad (6)$$

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$$v_{t1} = v_t \cos \phi, \quad v_{t2} = v_t \sin \phi \quad (6)$$

each R_f is used once

Number of pairs to be tested

$$N_{coll} = \frac{N_p^2 F_N}{2V_C} (\sigma_t g_r)_{max} \Delta t \quad (7)$$

N_p number of model particles in cell

F_N is representation, number of real particles represented by a model one.

V_C is volume of cell

g_r relative speed of interacting particles

σ_t total cross section. function of g_r

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a randomly chosen pair is accepted if
$$\frac{\sigma_t g_r}{(\sigma_t g_r)_{max}} > R_f \quad (8)$$

faster particles collide more frequently

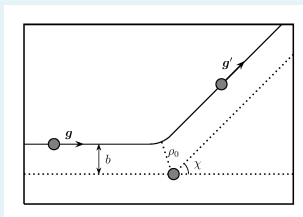
Once a pair is accepted

velocities v_1, v_2 are recalculated according to binary collision dynamics

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Scheme of binary collision

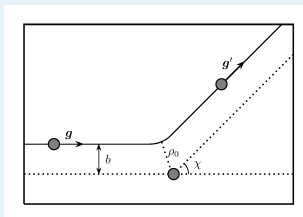


$g = g'$ relative speed does not change its magnitude

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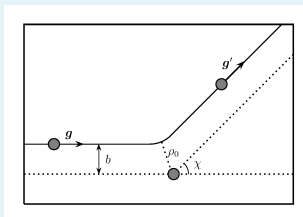
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deflection angle $0 \leq \chi \leq \pi$

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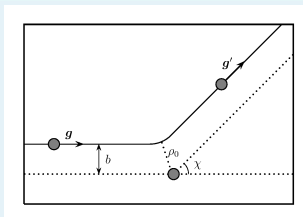
$g = g'$ relative speed does not change its magnitude

deflection angle $0 \leq \chi \leq \pi$ depends on potential

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velocities $\mathbf{v}_1, \mathbf{v}_2$ are recalculated according to binary collision dynamics

Scheme of binary collision



$g = g'$ relative speed does not change its magnitude

deflection angle $0 \leq \chi \leq \pi$ depends on potential

azimuthal impact angle $0 \leq \varepsilon \leq 2\pi$ determines the collision plane

$$\mathbf{v}, \mathbf{v}_* \Rightarrow \chi, \varepsilon \Rightarrow \mathbf{v}', \mathbf{v}'_*,$$

(9)

Hard sphere potential

Total cross section $\sigma_t = \pi d^2$ is constant.

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All directions of \mathbf{g}' are equiprobable

$$\cos \chi = 2R_f - 1, \quad \varepsilon = 2\pi R_f \quad (10)$$

each R_f is used once

Hard sphere potential

Total cross section $\sigma_t = \pi d^2$ is constant.

All directions of \mathbf{g}' are equiprobable

$$\cos \chi = 2R_f - 1, \quad \varepsilon = 2\pi R_f \quad (10)$$

each R_f is used once It is simple, but it leads to a wrong viscosity

$$\mu \propto T^{1/2} \quad (11)$$

Critics of variable hard sphere

All directions of \mathbf{g}' are equiprobable.

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All directions of \mathbf{g}' are equiprobable.

Total cross section $\sigma_t = \pi d^2$ depends on g .

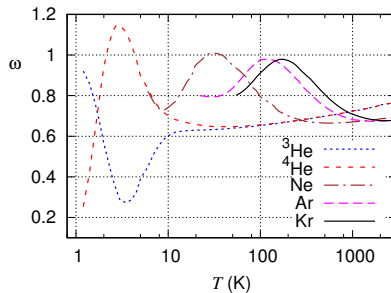
Diameter d is calculated assuming

$$\mu = \mu_{ref} \left(\frac{T}{T_{ref}} \right)^\omega \quad (12)$$

index ω is constant

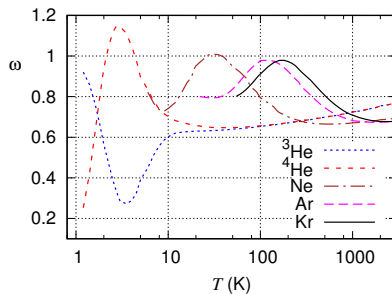
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Critics of Variable hard sphere

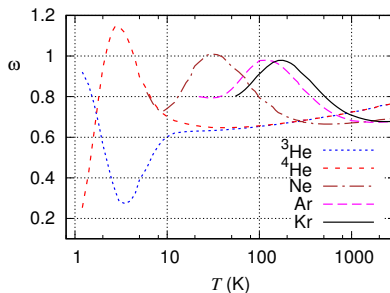
ω is not constant



Parameters d and ω are still fitting

Critics of Variable hard sphere

ω is not constant



Parameters d and ω are still fitting

It is not clear ω for collisions between different species.

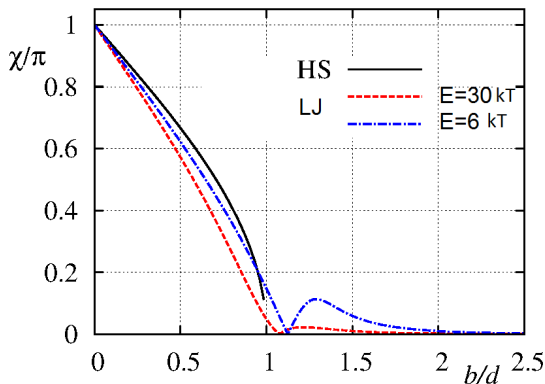
Arbitrary potential. Classical approach (Newton's laws)

$$\chi = \chi(b, g)$$

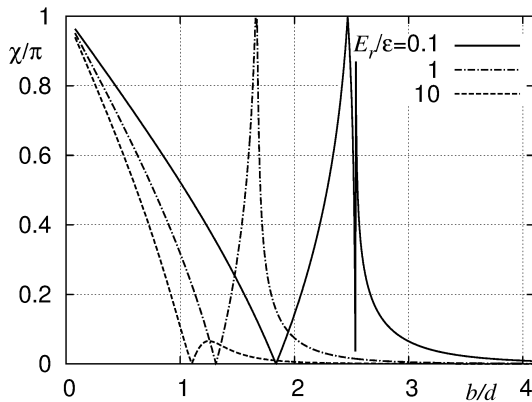
Arbitrary potential. Classical approach (Newton's laws)

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Deflection angle vs. impact parameter. Lennard-Jones potential.



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Arbitrary potential. Quantum approach

Schrödinger equation is applied

$\chi(g)$ is determined by differential cross section $\sigma(\chi, g)$

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Arbitrary potential

Calculation of deflection angle χ in each collision takes a long time by both classical and quantum approaches.

Arbitrary potential. Quantum approach

Schrödinger equation is applied

$\chi(g)$ is determined by differential cross section $\sigma(\chi, g)$

Arbitrary potential

Calculation of deflection angle χ in each collision takes a long time by both classical and quantum approaches.

χ are calculated for many values of g and stored in lookup tables.

Sharipov & Strapasson, *Phys. Fluids* **24** (2012)

Arbitrary potential. Lookup tables.

Matrix χ_{ij} is precalculated for discrete values of g_j .

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For a fixed j , all χ_{ij} are distributed so that all of them are equiprobable so that

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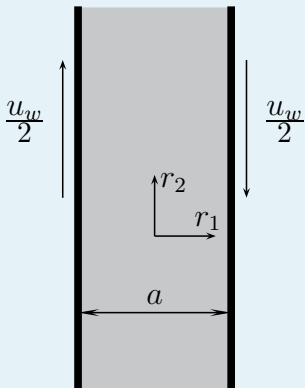
χ_{ij} are given in Supplementary material of several papers.

Computational effort with lookup tables is the same as for HS.

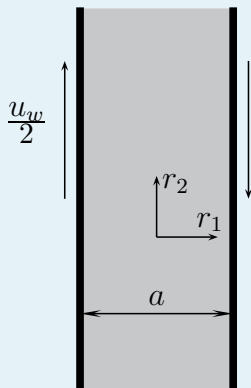
Hard sphere vs. *ab initio* potential

Classical approach vs. Quantum approach

Couette flow. Sharipov & Strapasson, *Phys. Fluids* **25** 027101 (2013)

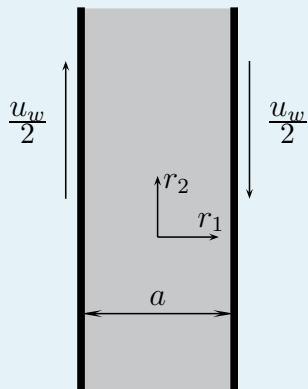


Couette flow. Sharipov & Strapasson, *Phys. Fluids* **25** 027101 (2013)



Shear stress P_{12} is calculated.

$$\Pi = -\frac{P_{12}}{p_0} \frac{v_0}{u_w}$$

Couette flow. Sharipov & Strapasson, *Phys. Fluids* **25** 027101 (2013)

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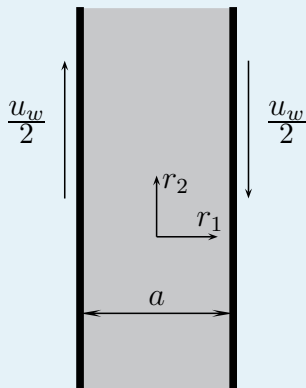
$$\Pi = -\frac{P_{12}}{p_0} \frac{v_0}{u_w} \quad (13)$$

as a function of mole fraction and rarefaction parameter

$$C = \frac{n_1}{n_1 + n_2}, \quad \delta = \frac{a p_0}{\mu v_0} \quad (14)$$

$$v_0 = \sqrt{\frac{2kT_0}{m}}, \quad m = C m_1 + (1 - C_1) m_2 \quad (15)$$

Couette flow. Sharipov & Strapasson, *Phys. Fluids* **25** 027101 (2013)



Numerical error 0.5%

Number of cells 400

Number of particles 40 000

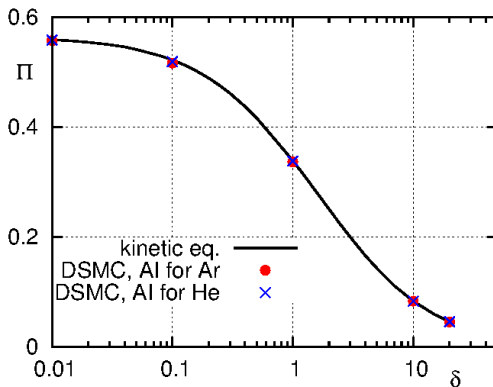
Time step $0.002a/v_0$

Number of steps 10^6 for $u_w/v_0 = 2$

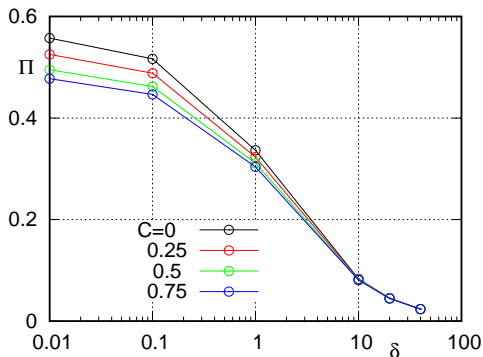
Number of steps 10^7 for $u_w/v_0 = 0.2$

Couette flow, Comparison with Discrete velocity method applied to BGK,

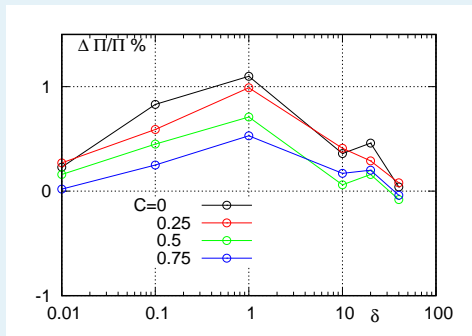
$$u_w = 0.2v_0$$



Numerical results based on AI potential, He-Ar, $u_w/v_0 = 0.2$.
 Sharipov & Strapasson, *Phys. Fluids* **25** 027101 (2013)

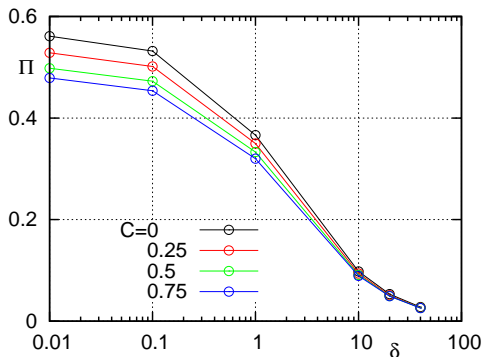


Relative difference between results based on AI and HS, $u_w/v_0 = 0.2$

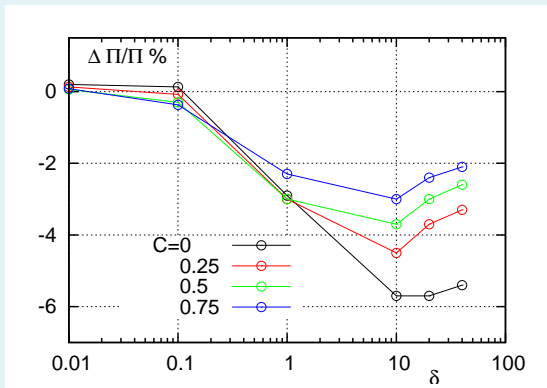


$$\frac{\Delta \Pi}{\Pi} = \frac{\Pi^{(HS)} - \Pi^{(AI)}}{\Pi^{(AI)}} \times 100\% \quad (16)$$

Numerical results based on AI potential, He-Ar, $u_w/v_0 = 2$.
 Sharipov & Strapasson, *Phys. Fluids* **25** 027101 (2013)

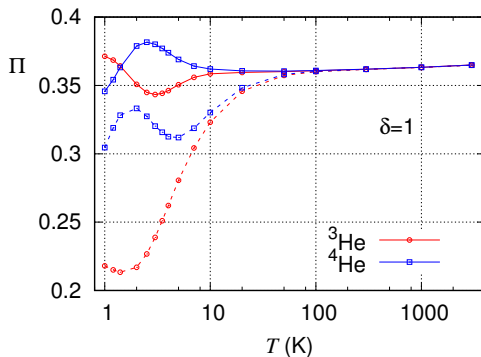


Relative difference between results based on AI and HS, $u_w/v_0 = 2$



$$\frac{\Delta \Pi}{\Pi} = \frac{\Pi^{(HS)} - \Pi^{(AI)}}{\Pi^{(AI)}} \times 100\% \quad (17)$$

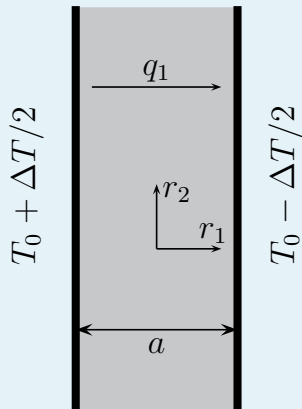
Classical and quantum approaches, $u_w/v_0 = 2$.
 Sharipov, *Physica A* **508**, 797-805 (2018)



Solid line - quantum approach
 dashed line - classical approach

Heat transfer between two plates.

Sharipov & Strapasson, *Int. J. Heat Mass Transfer* **71**, 91-97 (2014)



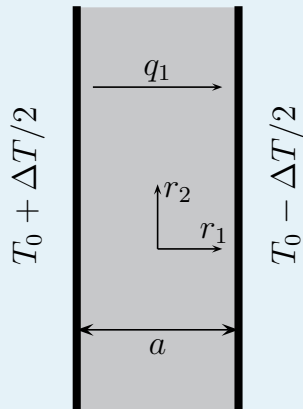
Heat q_x flow is calculated

$$Q = -\frac{q_x}{p_0 v_0} \frac{T_0}{\Delta T} \quad (18)$$

as function of C and δ

Heat transfer.

Sharipov & Strapasson, *Int. J. Heat Mass Transfer* **71**, 91-97 (2014)



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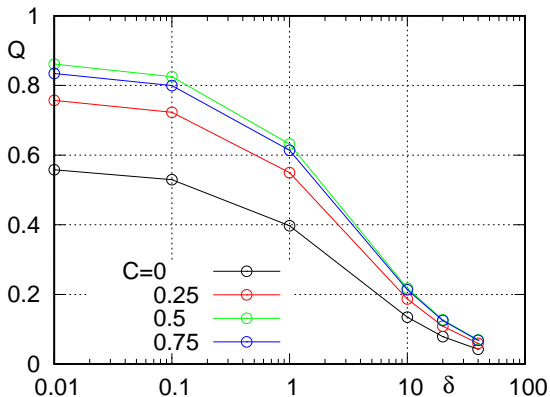
Number of particles 40 000

Time step $0.002a/v_0$

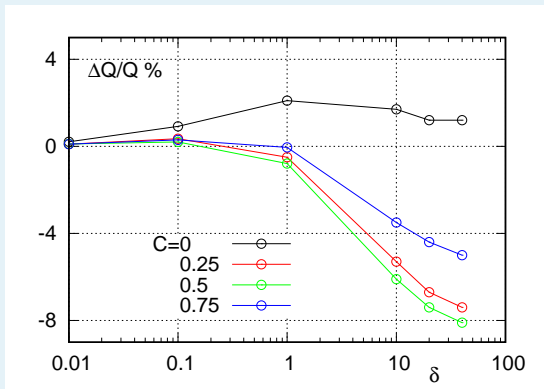
Number of steps 5×10^6 for $\Delta T/T_0 = 1.5$

Number of steps 10^7 for $\Delta T/T_0 = 0.2$

Numerical results based on AI potential, He-Ar, $\Delta T/T_0 = 0.2$.
 Sharipov & Strapasson, *Int. J. Heat Mass Transfer* **71**, 91-97 (2014)

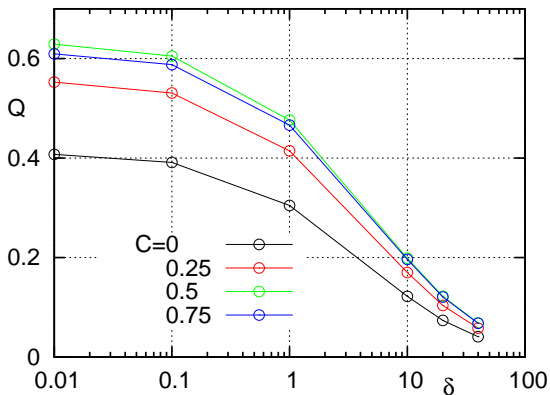


Relative difference between results based on AI and HS, $\Delta T/T_0 = 0.2$

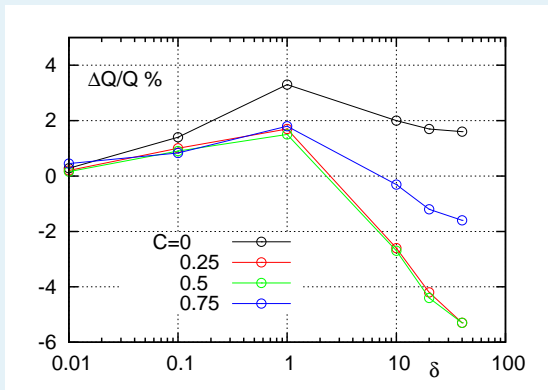


$$\frac{\Delta Q}{Q} = \frac{Q^{(AI)} - Q^{(HS)}}{Q^{(AI)}} \times 100\% \quad (19)$$

Numerical results based on Al potential, He-Ar, $\Delta T/T_0 = 1.5$.
 Sharipov & Strapasson, *Int. J. Heat Mass Transfer* **71**, 91-97 (2014)

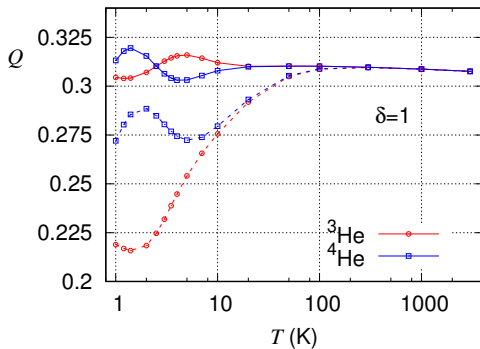


Relative difference between results based on AI and HS, $\Delta T/T_0 = 1.5$

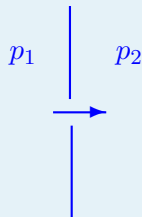
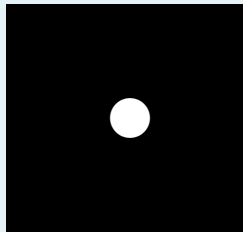


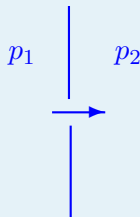
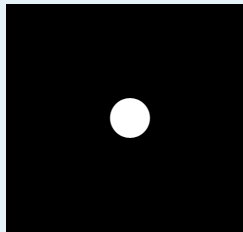
$$\frac{\Delta Q}{Q} = \frac{Q^{(AI)} - Q^{(HS)}}{Q^{(AI)}} \times 100\% \quad (20)$$

Classical and quantum approaches, $\Delta T/T_0 = 1.5$.
 Sharipov, *Physica A* **508**, 797-805 (2018)



Solid line - quantum approach
 dashed line - classical approach

Orifice flow. Sharipov & Strapasson, *Vacuum* **109**, 246-252 (2014)

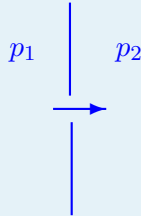
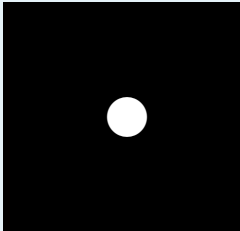
Orifice flow. Sharipov & Strapasson, *Vacuum* **109**, 246-252 (2014)

Reduced flow rate

$$W = \frac{\dot{M}}{\dot{M}_0}$$

$$\dot{M}_0 = \sqrt{\pi} a^2 p_1 / v_m \quad \text{free-molecular flow into vacuum}$$

Orifice flow. Sharipov & Strapasson, *Vacuum* 109, 246-252 (2014)

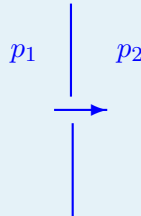
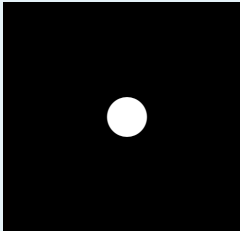


Determining parameters

➡ Rarefaction parameter

$$\delta = \frac{ap_1}{\mu v_m}$$

Orifice flow. Sharipov & Strapasson, *Vacuum* 109, 246-252 (2014)



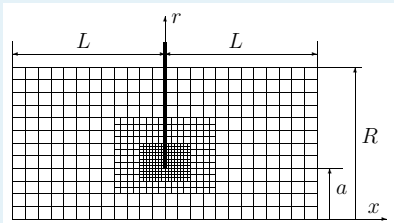
Determining parameters

➡ Rarefaction parameter

$$\delta = \frac{ap_1}{\mu v_m}$$

➡ Pressure ratio p_2/p_1

DSMC scheme



Sharipov & Strapasson, *Vacuum* **109** (2014)

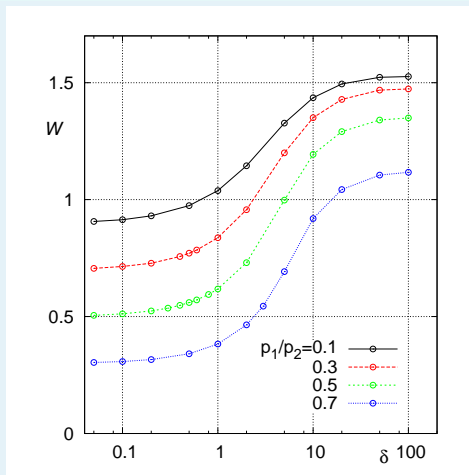
AI potential for He, Ar, Kr

Number of model particles - 3×10^7

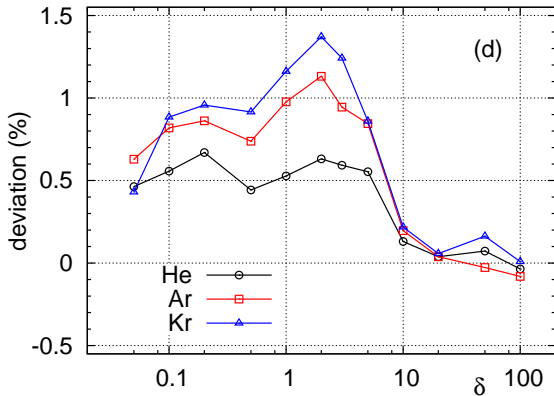
Time step - $0.005R/v_m$

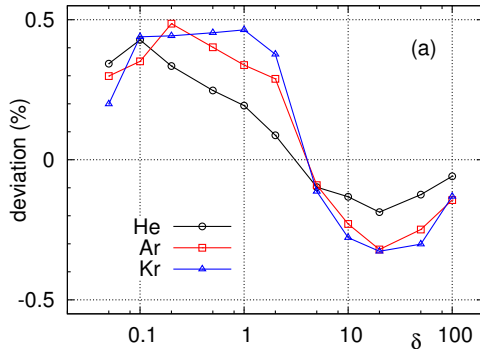
Number of sample - 10^6

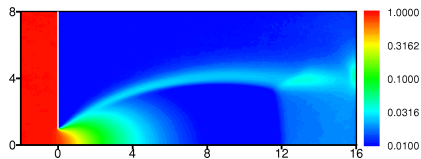
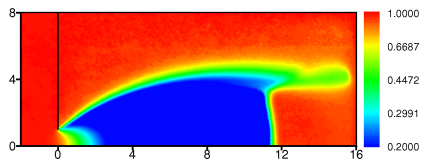
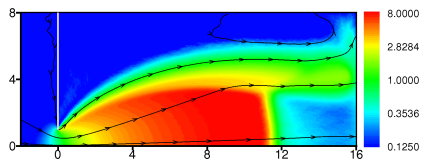
Numerical results based on AI potential, Ar.



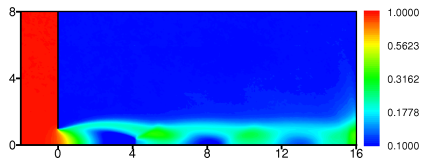
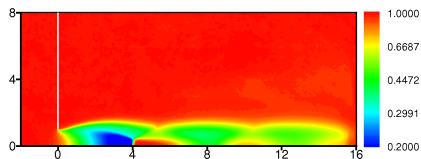
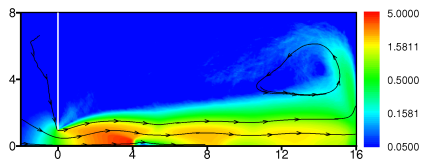
Sharipov & Strapasson, *Vacuum* **109** (2014)

Relative difference of flow rate between AI and HS, $p_1/p_0 = 0.7$ 

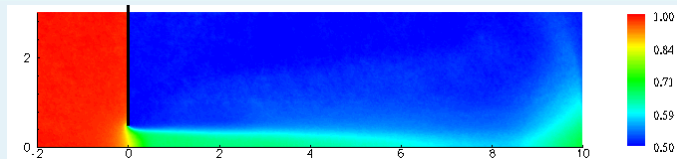
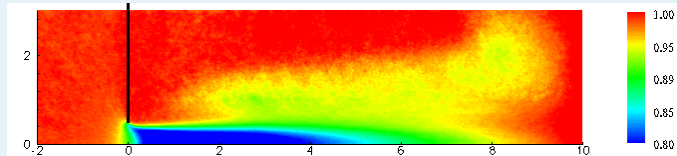
Relative difference of flow rate between AI and HS, $p_1/p_0 = 0.1$ Sharipov & Strapasson, *Vacuum* **109** (2014)

Flow-field at $p_1/p_0 = 0.01$ and $\delta = 1000$  ρ/ρ_0 density T/T_0 temperature

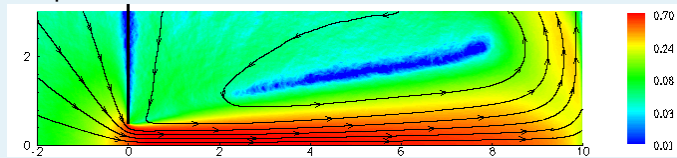
Local Mach number

Flow-field at $p_1/p_0 = 0.1$ and $\delta = 1000$  ρ/ρ_0 density T/T_0 temperature

Mach number

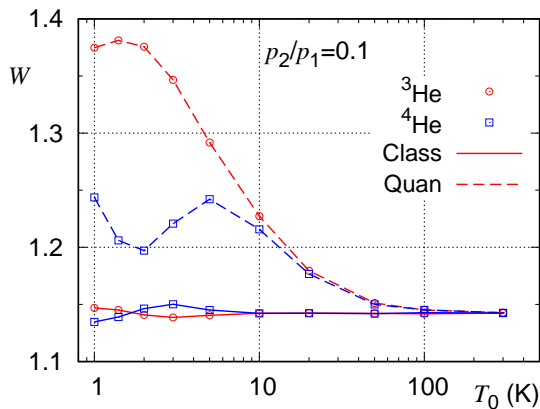
Flow-field at $p_1/p_0 = 0.5$ and $\delta = 1000$  ρ/ρ_0 density T/T_0

temperature



Mach number

AI potential, quantum effects.

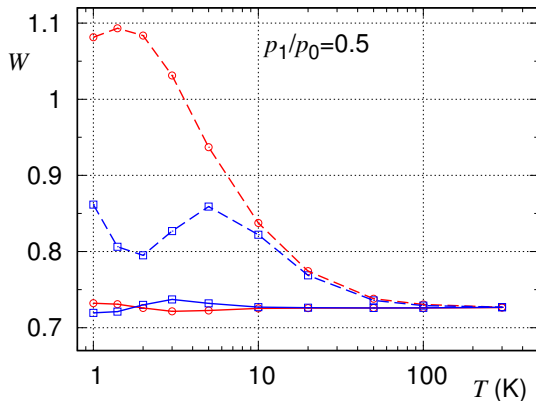


$$\delta = 2$$

solid - quantum,
dashed - classical

Sharipov, *Vacuum* **156**. 146 (2018)

AI potential, quantum effects.

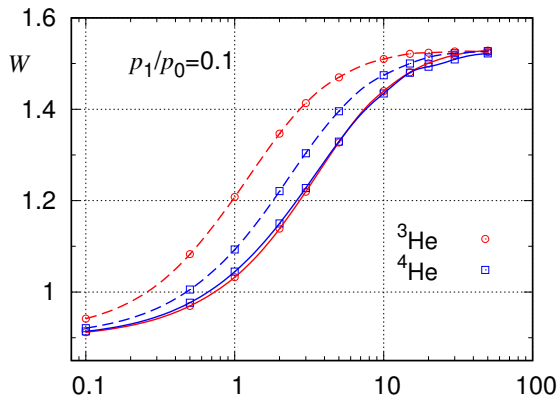


$$\delta = 2$$

solid - quantum,
dashed - classical

Sharipov, *Vacuum* **156**. 146 (2018)

AI potential, quantum effects.



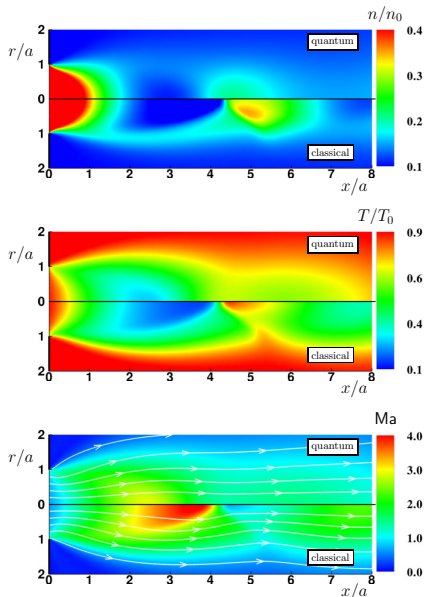
$T = 3\text{ K}$

solid - quantum,

dashed - classical

Sharipov, *Vacuum* **156**. 146 (2018)

AI potential, quantum effects



$$p_1/p_0 = 0.1,$$

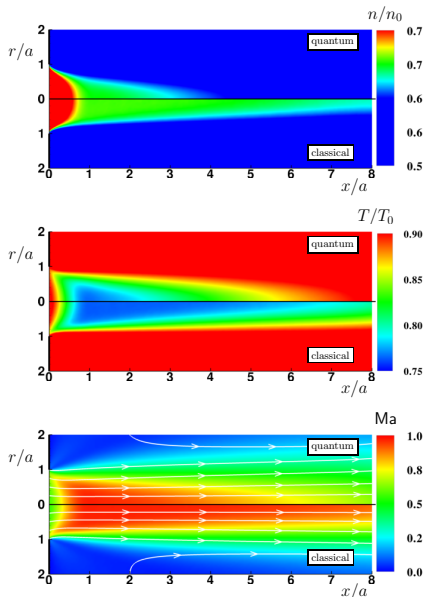
$$\delta = 50$$

$$T = 3 \text{ K}$$

cl. CPU time = $4 \times$

quan. CPU time

AI potential, quantum effects



$$p_1/p_0 = 0.5,$$

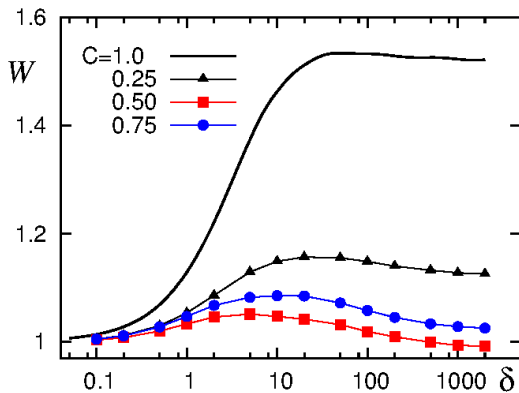
$$\delta = 50$$

$$T = 3 \text{ K}$$

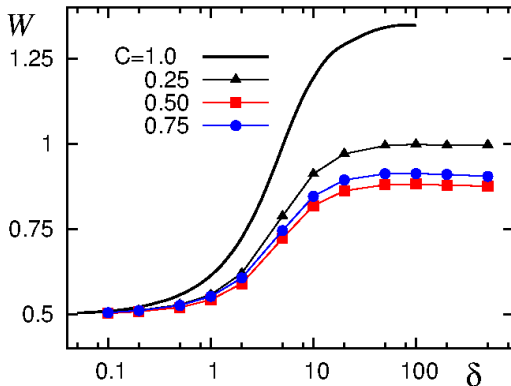
cl. CPU time = $4 \times$

quan. CPU time

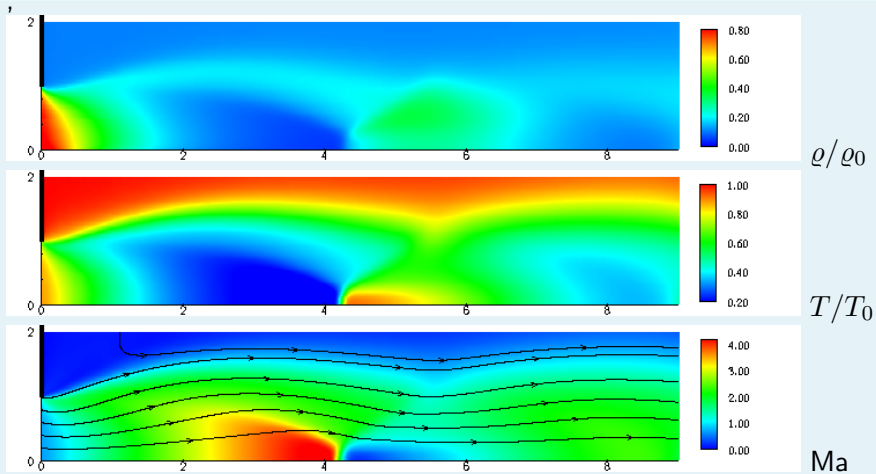
Reduced flow rate W vs δ at $p_1/p_0 = 0$, mixture of He-Ar.
 Sharipov, *Vacuum* **143** (2017)



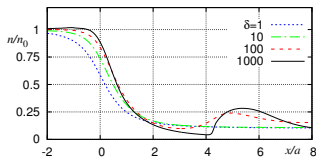
Reduced flow rate W vs δ at $p_1/p_0 = 0.5$, mixture of He-Ar.
 Sharipov, *Vacuum* **143** (2017)



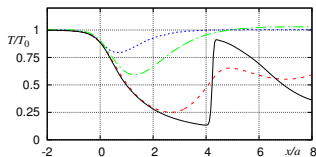
Flow-field of equimolar mixture He-Ar at $p_1/p_0 = 0.1$ and $\delta = 1000$, Sharipov, *Vacuum* **143** (2017)



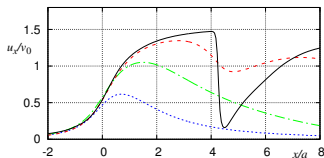
Axial distribution of equimolar mixture at $p_1/p_0 = 0.1$, Sharipov, *Vacuum* 143 (2017)



density

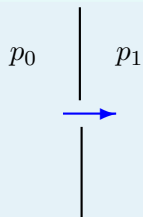
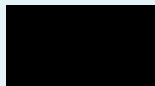


temperature

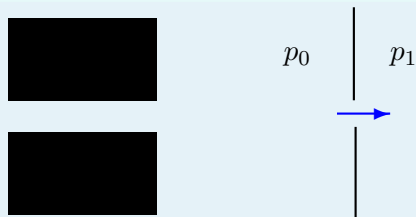


bulk velocity

Scheme of slit flow



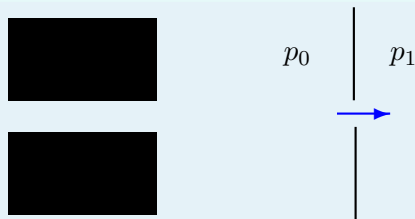
Scheme of slit flow



Reduced flow rate

$$W = \dot{M} / \dot{M}_0, \quad \dot{M}_0 = p_0 H / \sqrt{\pi} v_m, \quad v_m = \sqrt{2kT/m} \quad (21)$$

Scheme of slit flow



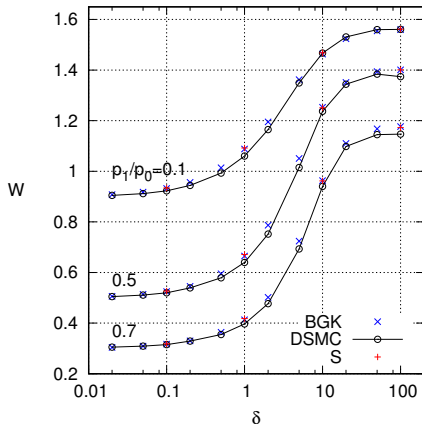
Reduced flow rate

$$W = \dot{M}/\dot{M}_0, \quad \dot{M}_0 = p_0 H / \sqrt{\pi} v_m, \quad v_m = \sqrt{2kT/m} \quad (21)$$

Rarefaction parameter

$$\delta = p_0 H / \mu v_m \propto 1/\text{Kn} \quad (22)$$

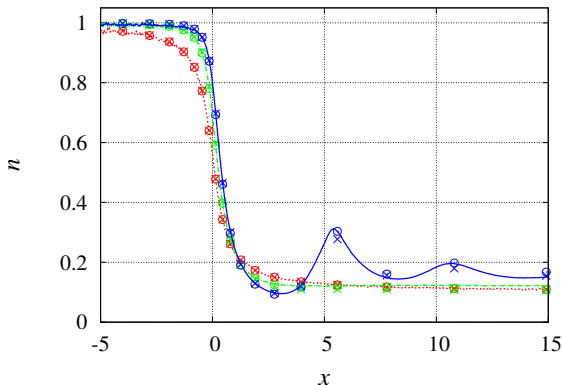
Flow rate through a slit. BGK vs. DSMC



$$\Delta W/W < 4\%$$

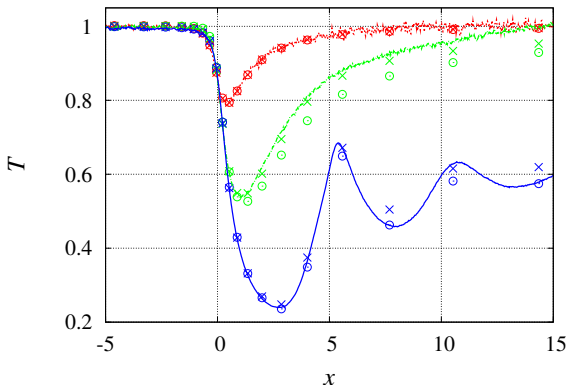
DSMC: Sharipov & Kozak, *Eur. J. Mech. B/Fluids* **30**(5), (2011)

BGK and S: Graur, Polikarpov & Sharipov, *ZAMP* **63**, 503-50 (2012)

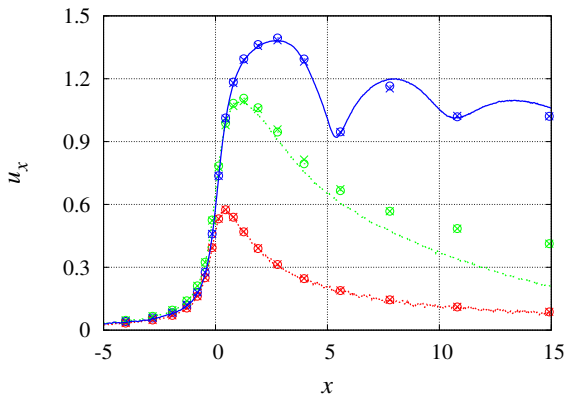
Axial distribution of flowfield past a slit at $p_1/p_0 = 0.1$. BGK vs. DSMC

Density, Lines - DSMC, circles - BGK, crosses - S-model;

red - $\delta = 0.1$, green - $\delta = 10$, blue - $\delta = 100$

Axial distribution of flowfield past a slit at $p_1/p_0 = 0.1$. BGK vs. DSMC

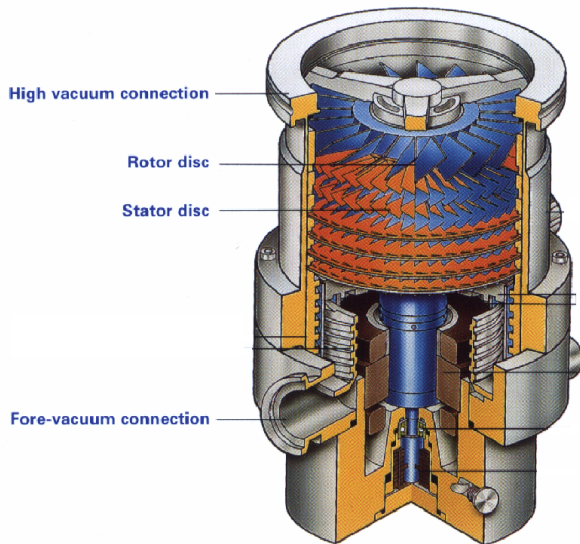
Temperature, Lines - DSMC, circles - BGK, crosses - S-model;
 red - $\delta = 0.1$, green - $\delta = 10$, blue - $\delta = 100$

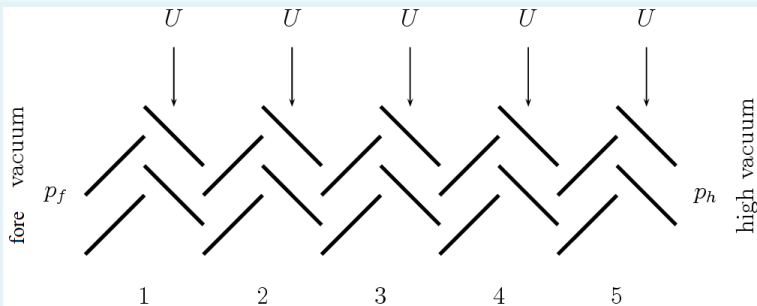
Axial distribution of flowfield past a slit at $p_1/p_0 = 0.1$. BGK vs. DSMC

Bulk velocity, Lines - DSMC, circles - BGK, crosses - S-model;
 red - $\delta = 0.1$, green - $\delta = 10$, blue - $\delta = 100$

Modelling of vacuum pumps

Scheme of pump



Scheme of pump, Sharipov, *J. Vac. Sci. Technol. A* **28**, 1312 (2010)

Pumping speed

$$S = \frac{2\sqrt{\pi}}{v_m A} \frac{q}{p_f} \quad (23)$$

Pumping speed

$$S = \frac{2\sqrt{\pi}}{v_m A} \frac{q}{p_f} \quad (23)$$

A - working area of TMP

v_m - the most probable speed of molecules

q - throughput [Pa m³/s]

p_f - fore vacuum pressure

Compression ratio

$$k = \frac{\text{fore vacuum pressure}}{\text{high vacuum pressure}} \quad (24)$$

Main determining parameters

- $$U = \frac{\text{blade speed}}{\text{most probable molecular speed}} \quad (25)$$

Main determining parameters



$$U = \frac{\text{blade speed}}{\text{most probable molecular speed}} \quad (25)$$

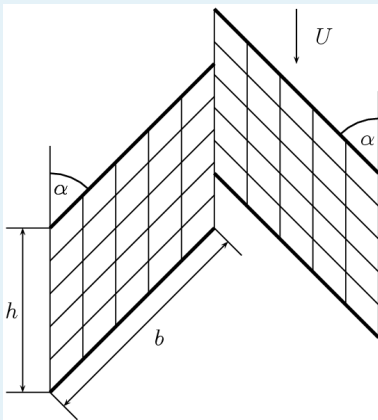
- Rarefaction parameter

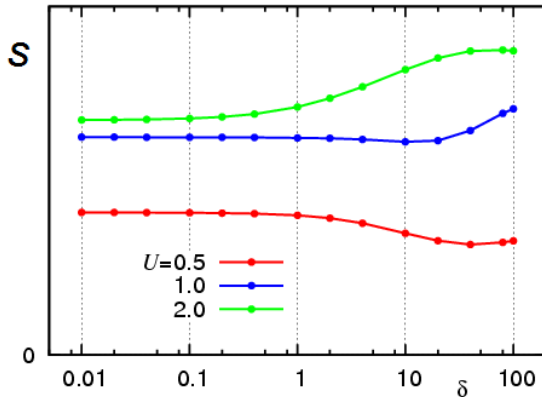
$$\delta = \frac{hp_f}{\mu v_m} \quad (26)$$

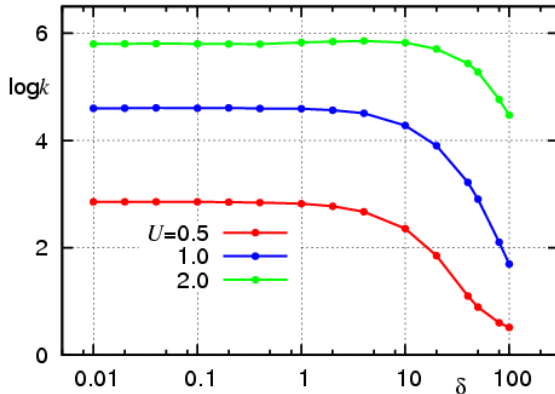
h space between blades

μ gas viscosity

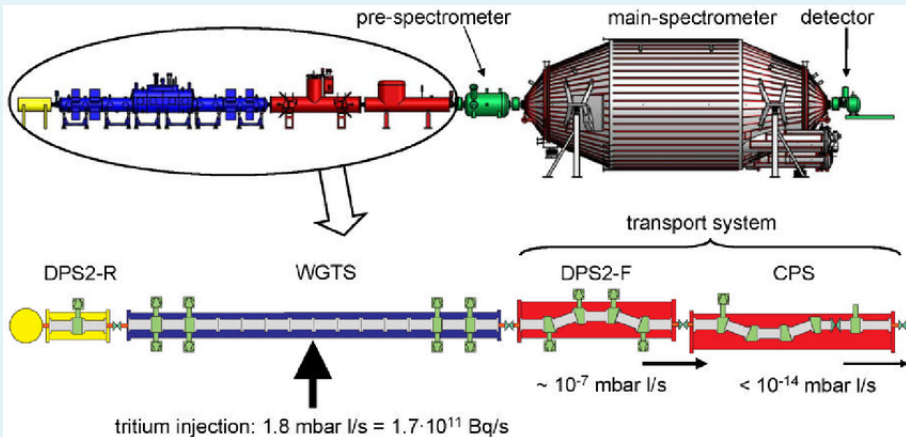
Numerical scheme, DSMC



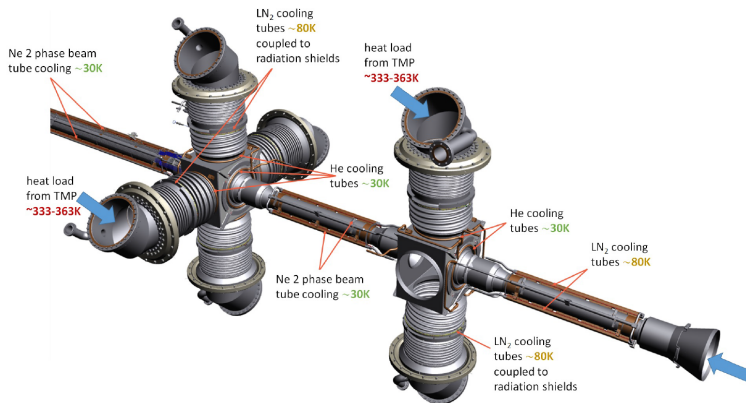
Results, Maximum pumping speed, $p_h = p_f$ 

Results, Limit compression ratio, $S = 0$ 

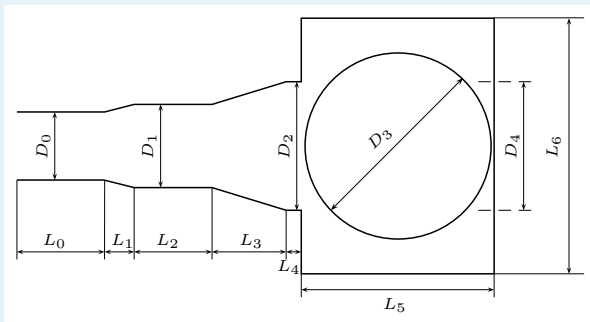
Measurement of neutrino mass (KATRIN)



Scheme of flow

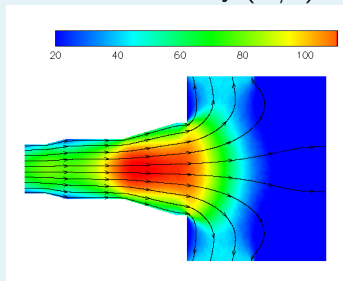


Scheme of flow

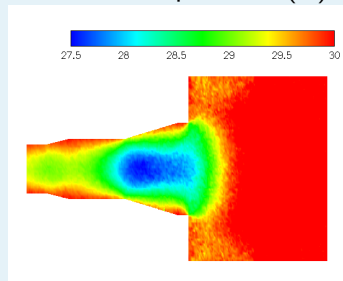


Flow-field. All pumps work

bulk velocity (m/s)

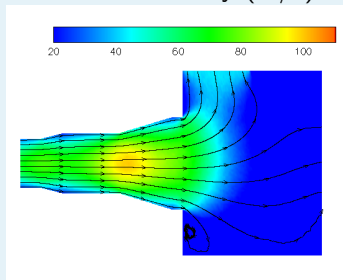


temperature (K)

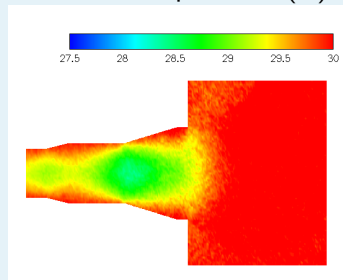


Flow-field. One pump fails

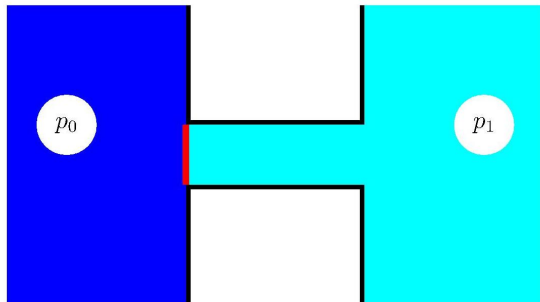
bulk velocity (m/s)



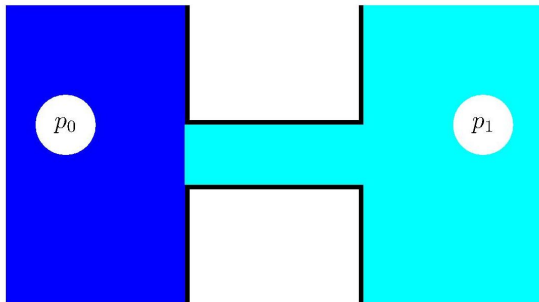
temperature (K)



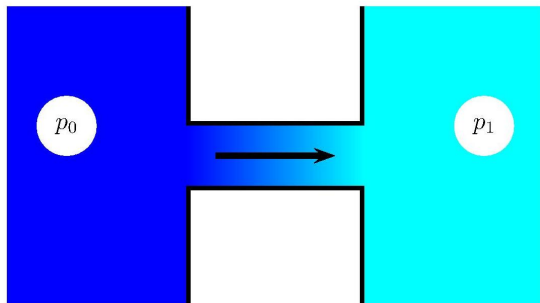
Transient flows through short tube



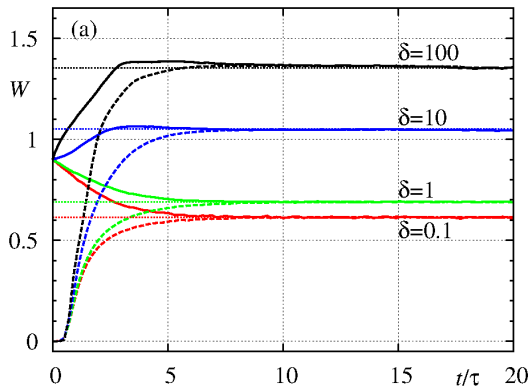
Transient flows through short tube



Transient flows through short tube

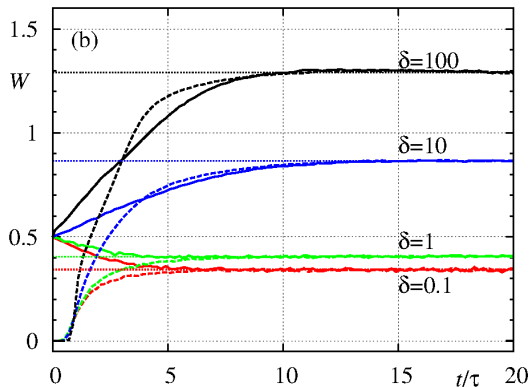


Flow rate vs time t ($\tau = a/v_m$) at $L/a = 1$ and $p_2/p_1 = 0.1$
 solid - inlet, dashed - outlet



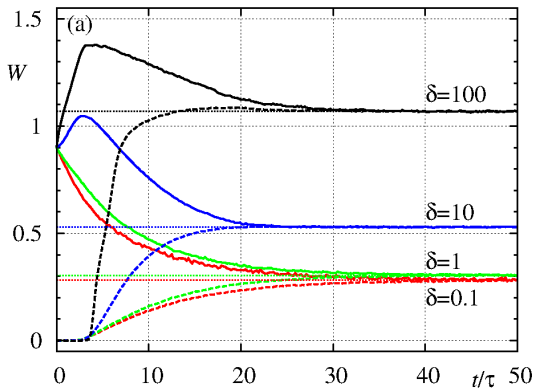
Sharipov, *Vacuum* **90**, P.25 (2013)

Flow rate vs time t ($\tau = a/v_m$) at $L/a = 1$ and $p_2/p_1 = 0.5$
 solid - inlet, dashed - outlet



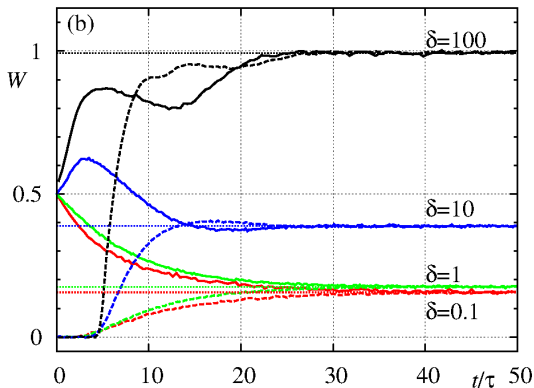
Sharipov, *Vacuum* **90**, P.25 (2013)

Flow rate vs time t ($\tau = a/v_m$) at $L/a = 5$ and $p_2/p_1 = 0.1$
 solid - inlet, dashed - outlet



Sharipov, *Vacuum* **90**, P.25 (2013)

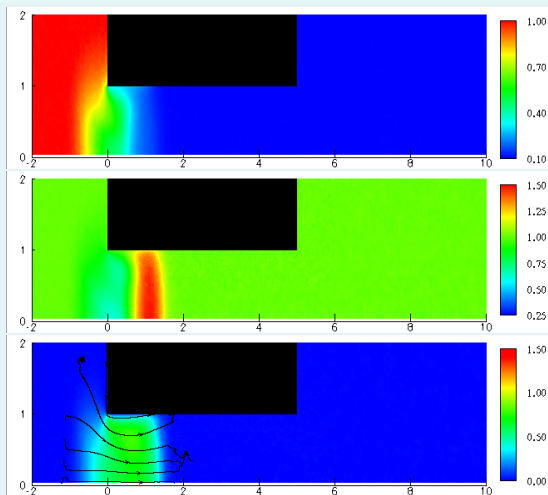
Flow rate vs time t ($\tau = a/v_m$) at $L/a = 5$ and $p_2/p_1 = 0.5$
 solid - inlet, dashed - outlet



Sharipov, *Vacuum* **90**, P.25 (2013)

Flowfield $p_1/p_0 = 0.1$, $L/a = 5$, $\delta = 100$

$$t = 1 \times \tau, \quad \tau = a/v_m$$



Density

$$\rho/\rho_0$$

Temperature

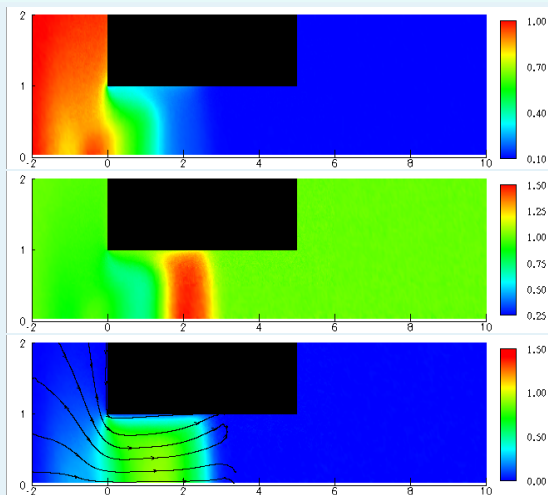
$$T/T_0$$

Bulk velocity

$$u_x/v_m$$

Flowfield $p_1/p_0 = 0.1$, $L/a = 5$, $\delta = 100$

$$t = 2 \times \tau, \quad \tau = a/v_m$$



Density

$$\rho/\rho_0$$

Temperature

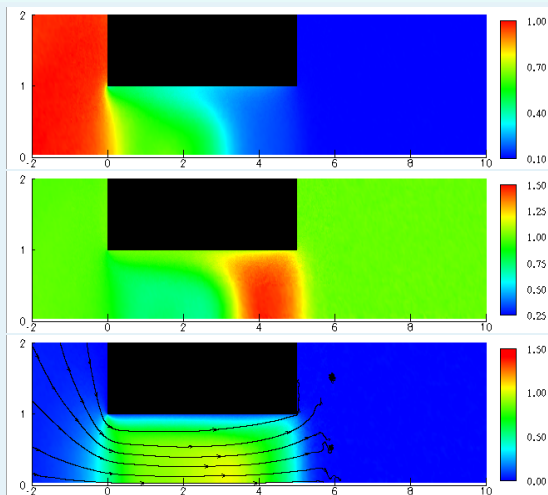
$$T/T_0$$

Bulk velocity

$$u_x/v_m$$

Flowfield $p_1/p_0 = 0.1$, $L/a = 5$, $\delta = 100$

$$t = 4 \times \tau, \quad \tau = a/v_m$$



Density

ρ/ρ_0

Temperature

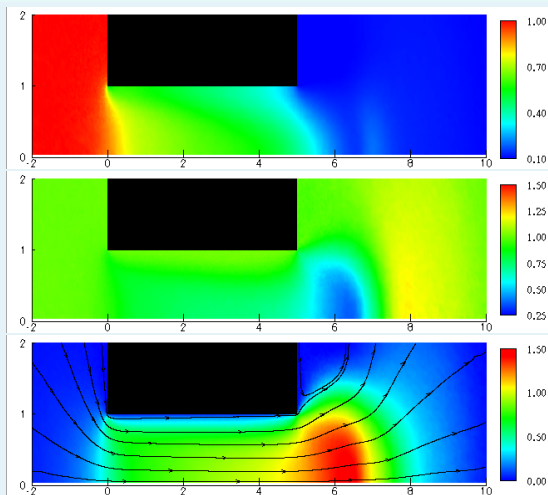
T/T_0

Bulk velocity

u_x/v_m

Flowfield $p_1/p_0 = 0.1$, $L/a = 5$, $\delta = 100$

$$t = 8 \times \tau, \quad \tau = a/v_m$$



Density

$$\rho/\rho_0$$

Temperature

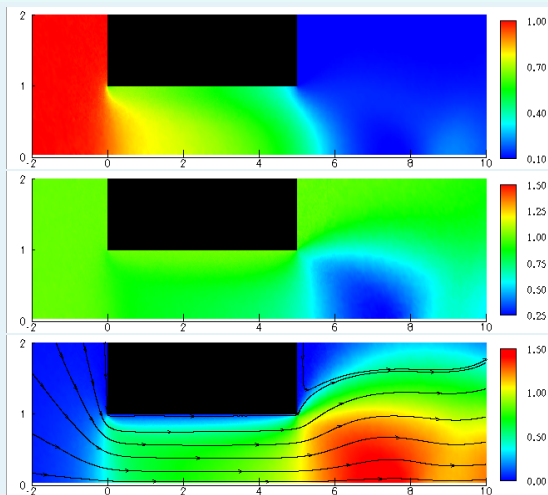
$$T/T_0$$

Bulk velocity

$$u_x/v_m$$

Flowfield $p_1/p_0 = 0.1$, $L/a = 5$, $\delta = 100$

$$t = 15 \times \tau, \quad \tau = a/v_m$$



Density

$$\rho/\rho_0$$

Temperature

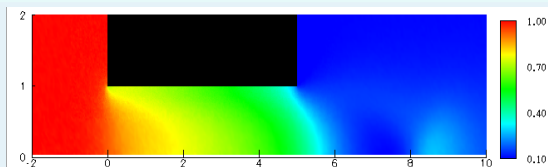
$$T/T_0$$

Bulk velocity

$$u_x/v_m$$

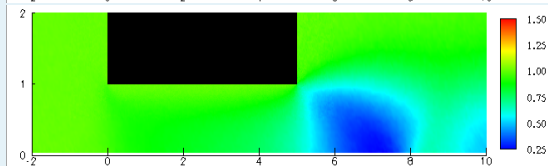
Flowfield $p_1/p_0 = 0.1$, $L/a = 5$, $\delta = 100$

$$t = 30 \times \tau, \quad \tau = a/v_m$$



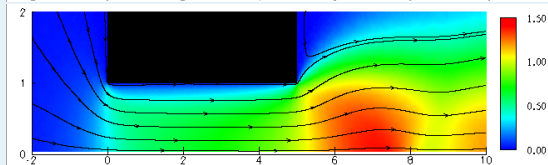
Density

$$\rho/\rho_0$$



Temperature

$$T/T_0$$

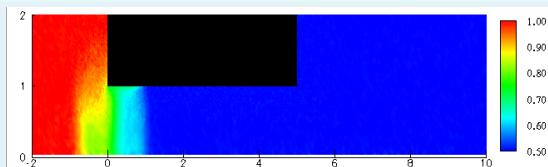


Bulk velocity

$$u_x/v_m$$

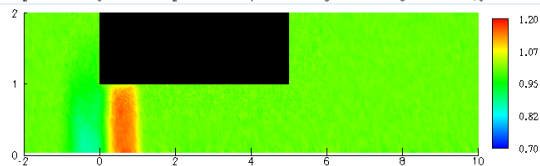
Flowfield $p_1/p_0 = 0.5$, $L/a = 5$, $\delta = 100$

$$t = 1 \times \tau, \quad \tau = a/v_m$$



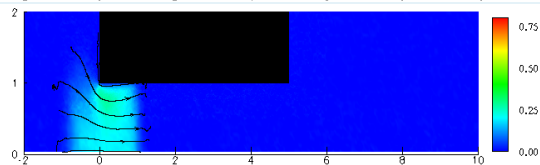
Density

$$\rho/\rho_0$$



Temperature

$$T/T_0$$

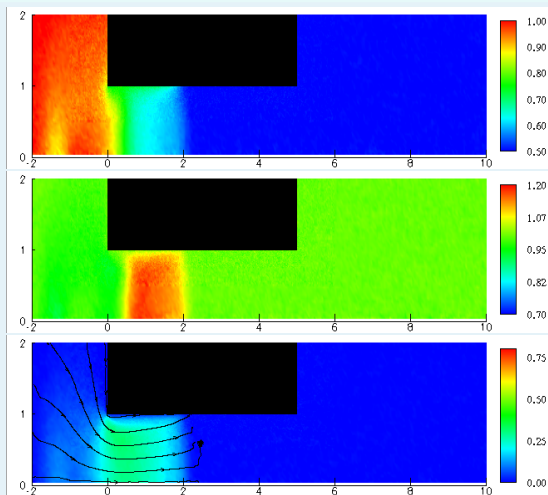


Bulk velocity

$$u_x/v_m$$

Flowfield $p_1/p_0 = 0.5$, $L/a = 5$, $\delta = 100$

$$t = 2 \times \tau, \quad \tau = a/v_m$$



Density

 ρ/ρ_0

Temperature

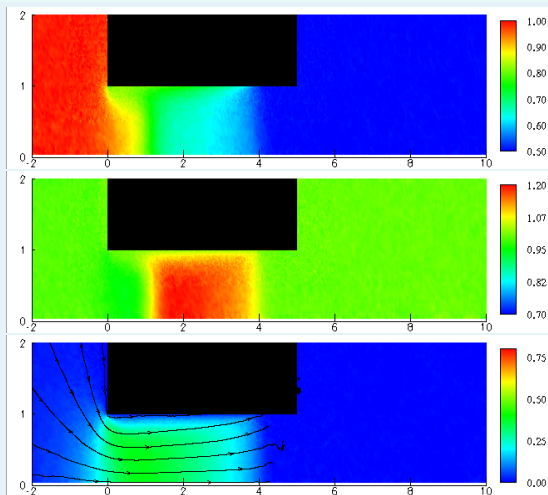
 T/T_0

Bulk velocity

 u_x/v_m

Flowfield $p_1/p_0 = 0.5$, $L/a = 5$, $\delta = 100$

$$t = 4 \times \tau, \quad \tau = a/v_m$$



Density

$$\rho/\rho_0$$

Temperature

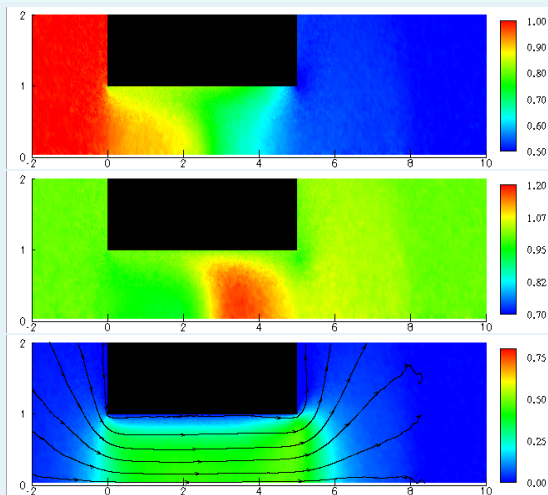
$$T/T_0$$

Bulk velocity

$$u_x/v_m$$

Flowfield $p_1/p_0 = 0.5$, $L/a = 5$, $\delta = 100$

$$t = 8 \times \tau, \quad \tau = a/v_m$$



Density

 ρ/ρ_0

Temperature

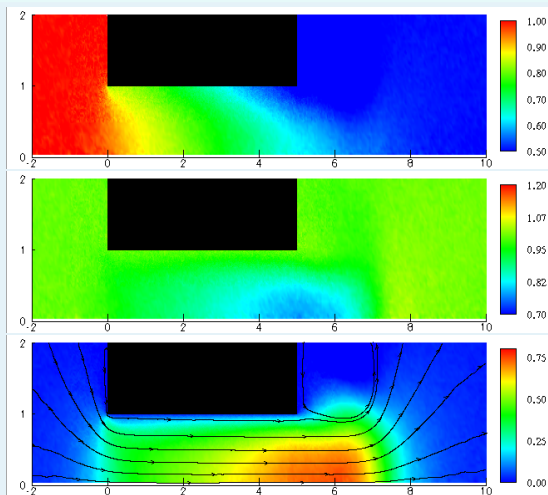
 T/T_0

Bulk velocity

 u_x/v_m

Flowfield $p_1/p_0 = 0.5$, $L/a = 5$, $\delta = 100$

$$t = 15 \times \tau, \quad \tau = a/v_m$$



Density

 ρ/ρ_0

Temperature

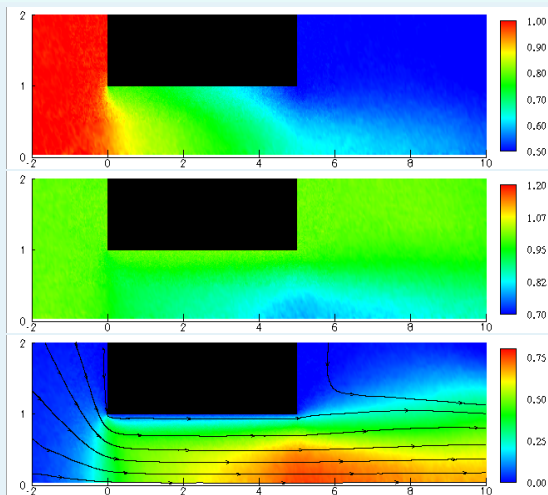
 T/T_0

Bulk velocity

 u_x/v_m

Flowfield $p_1/p_0 = 0.5$, $L/a = 5$, $\delta = 100$

$$t = 30 \times \tau, \quad \tau = a/v_m$$



Density

$$\rho/\rho_0$$

Temperature

$$T/T_0$$

Bulk velocity

$$u_x/v_m$$

TPMC, DSMC. Numerical codes

- G. A. Bird, *Molecular Gas Dynamics and the Direct Simulation of Gas Flows* (Oxford University Press, 1994)
- G. A. Bird, *The DSMC method* (2013)

Lookup tables of χ

F. Sharipov, Direct simulation Monte Carlo method based on ab initio potential: Recovery of transport coefficients of multi-component mixtures of noble gases. *Phys. Fluids* **34**, 097114 (2022).

THE END of Lecture 2

Thank you
for your attention
<http://fisica.ufpr.br/sharipov>